

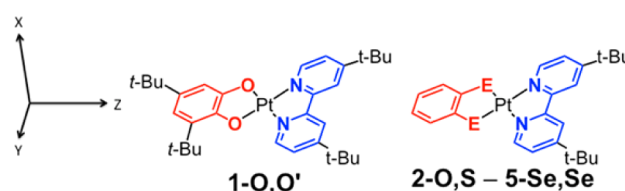
## Ligand Control of Donor–Acceptor Excited-State Lifetimes

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## Supporting Information

**ABSTRACT:** Transient absorption and emission spectroscopic studies on a series of diimineplatinum(II) dichalcogenolenes, LPTl', reveal charge-separated dichalcogenolene → diimine charge-transfer (LL'/CT) excited-state lifetimes that display a remarkable and nonperiodic dependence on the heteroatoms of the dichalcogenolene ligand. Namely, there is no linear relationship between the observed lifetimes and the principle quantum number of the E donors. The results are explained in terms of heteroatom-dependent singlet–triplet (S–T) energy gaps and anisotropic covalency contributions to the M–E (E = O, S, Se) bonding scheme that control rates of intersystem crossing. For the dioxolene complex, **1-O,O'**,  $E(T_2) > E(S_1)$  and rapid nonradiative decay occurs from  $S_1$  to  $S_0$ . However,  $E(T_2) \leq E(S_1)$  for the heavy-atom congeners, and this provides a mechanism for rapid intersystem crossing. Subsequent internal conversion to  $T_1$  in **3-S,S** produces a long-lived, emissive triplet. The two LPTl' complexes with mixed chalcogen donors and **5-Se,Se** show lifetimes intermediate between those of **1-O,O'** and **3-S,S**.

Square-planar compounds of the (dithiolene)Pt(diimine) series have generated considerable interest because of their rich photophysical properties and their remarkable photoluminescent behavior.<sup>1–8</sup> These d<sup>8</sup> platinum(II) compounds possess a flexible electronic structure that results in a low-energy ligand-to-ligand charge-transfer (LL'/CT) or mixed-metal ligand-to-ligand charge-transfer (MMLL'/CT) transition that absorbs strongly in the visible-to-near-infrared (NIR) region of the spectrum and possesses considerable dithiolene → diimine charge-transfer character. The LL'/CT transition in this class of compounds has been of particular interest to us in furthering our understanding of biradical charge-separated states because the LL'/CT excited state results in a large degree of charge separation along the long (z) axis of the molecule (Figure 1) to formally create an excited (dithiolene<sup>•+</sup>)Pt(diimine<sup>•-</sup>) LL'/CT state with open-shell donor–acceptor biradical character. Long-lived luminescence is observed following visible photoexcitation of the LL'/CT band for (bdt)Pt(di-*tert*-Bu-bpy), **3-S,S**, and other (dithiolene)Pt(diimine) complexes.<sup>2</sup> This emissive behavior is absent for the corresponding catechol complex, **1-O,O'**, which fully relaxes to the ground state within 630 ps.<sup>9</sup>



**Figure 1.** Structures of LPTl' complexes **1-O,O'**, **2-O,S**, **3-S,S**, **4-S,Se**, and **5-Se,Se**, where E,E = O,S, S,S, S,Se, and Se,Se. The donor ligand is depicted in red and the acceptor fragment in blue.

The ground-state electronic absorption spectra for **1-O,O'**, **2-O,S**, **3-S,S**, **4-S,Se**, and **5-Se,Se** are all characterized by the presence of a broad LL'/CT band with nearly equivalent absorption maxima that occur in a narrow energy range between 16950 cm<sup>-1</sup> for **1-O,O'** and 18350 cm<sup>-1</sup> for **3-S,S** (Table 1; for spectra, see the Supporting Information). The spectral similarities derive from the fact that they all possess similar  $a_2^2 b_1^2 b_1^0$  ground-state configurations (in effective  $C_{2v}$  symmetry), where the filled  $a_2^2$  and  $b_1^2$  donor orbitals are dominantly composed of – and + combinations of out-of-plane  $E_p$  orbitals, respectively. The  $b_1$  lowest unoccupied molecular orbital (LUMO) is primarily localized on the bipyridine acceptor. The electronic structure similarity among these compounds is further exemplified by the similar nature of their transient absorption (TA) difference spectra (Figure 2, inset). The TA difference spectra show a ground-state bleach of the LL'/CT band (545–564 nm), and the decay of the TA spectra is associated with a ground-state recovery time constant,  $\tau_{TA}$  (Table 1).

A Jablonski diagram describing the relevant photophysical processes for **1-O,O'**–**5-Se,Se** is presented in Figure 3. The diagram is composed of the  $S_0$  ground state, the  $S_1$  ( $a_2^2 b_1^1 b_1^1$ ;  $^1A_1$ ) and  $S_2$  ( $a_2^1 b_1^2 b_1^1$ ;  $^1B_2$ ) excited states, and their corresponding triplet  $T_1$  ( $^3A_1$ ) and  $T_2$  ( $^3B_2$ ) states. In the two-state limit (i.e., two-level approximation), the rate constant for intersystem crossing (ISC) to the triplet manifold is governed by spin–orbit coupling (SOC) according to<sup>10</sup>

$$k_{ISC} \propto \langle S_n | \mathbf{L}_x | T_n \rangle^2 / \Delta E_{ST}^2 \quad (1)$$

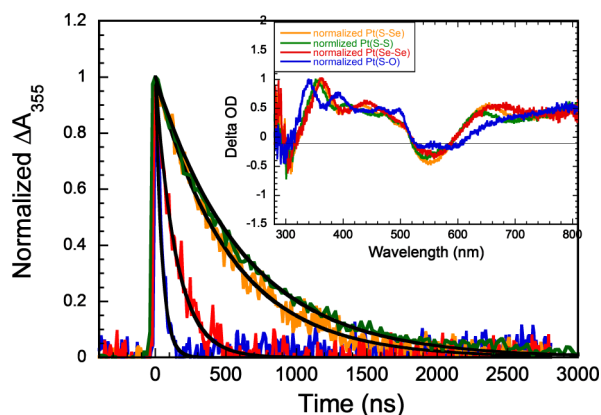
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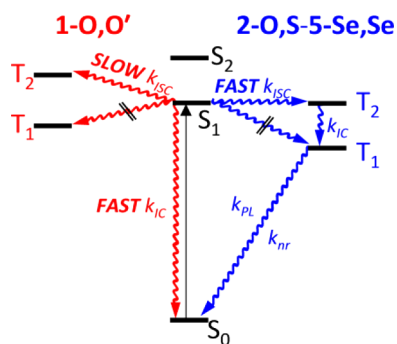


Table 1. Photophysical Properties of LL'Pt Complexes<sup>a</sup>

	1-O,O'	2-O,S	3-S,S	4-S,Se	5-Se,Se
$E_{\text{abs}} (\sim S_1-S_0)$ (cm <sup>-1</sup> )	16880 (592 nm)	17750 (564 nm)	18350 (545 nm)	18050 (554 nm)	18000 (555 nm)
$E_{\text{PL}}$ (cm <sup>-1</sup> )		13500 (740 nm)	14250 (702 nm)	14300 (700 nm)	14250 (702 nm)
Stokes (cm <sup>-1</sup> )		4217	4104	3765	3773
$\tau_{\text{PL}}$ (ns) <sup>a</sup>		46 ± 10	608 ± 10	540 ± 10	154 ± 10
$\tau_{\text{TA}}$ (ns) <sup>a</sup>	~0.6 <sup>b</sup>	23 ± 10	621 ± 10	511 ± 10	166 ± 10

<sup>a</sup>PL = photoluminescence; TA = transient absorption. <sup>b</sup>Reference 9.

**Figure 2.** (A) Kinetic traces for 2-O,S (blue), 3-S,S (green), 4-S,Se (orange), and 5-Se,Se (red) in degassed CH<sub>2</sub>Cl<sub>2</sub> at 298 K, obtained by monitoring compound-specific transient absorbances between 340 and 363 nm (excitation at the LL'/CT absorption maximum). Inset (B): TA difference spectra recorded immediately following a 5 ns gate delay (pump: 5 mJ/pulse at the LL'/CT absorption maximum).



**Figure 3.** Jablonski diagram for the observed photoprocesses in nonemissive 1-O,O' (left) and emissive complexes (right). Spin-orbit-promoted ISC from S<sub>1</sub> to T<sub>1</sub> is symmetry-forbidden, while spin-orbit-promoted ISC from S<sub>1</sub> to T<sub>2</sub> is symmetry-allowed. See the text.

Because  $\langle S_1 | L_z | T_1 \rangle = \langle A_1 | L_z | A_1 \rangle = 0$  ( $L_z$  are the orbital angular momentum operators, which transform as  $a_2$ ,  $b_1$ , and  $b_2$  in  $C_{2v}$  symmetry), the  $S_1 \rightarrow T_1$  process is symmetry-forbidden and  $k_{\text{ISC}}$  is expected to be small. The same holds true for  $T_1 \rightarrow S_0$  ISC, leading to ground-state recovery from the triplet manifold. In contrast, ISC between  $S_1$  and  $T_2$  is symmetry-allowed because  $\langle A_1 | L_z | B_2 \rangle \neq 0$ , and the ISC rate will be maximal when  $S_1$  and  $T_2$  are close in energy (eq 1).

Across the entire 1-O,O'–5-Se,Se series, the observed rates of ground-state recovery vary by nearly 3 orders of magnitude. The fast ground-state recovery rate for 1-O,O' (630 ps) occurs by  $S_1 \rightarrow S_0$  back-electron transfer (i.e., internal conversion, IC). Because the  $S_1 \rightarrow T_1$  ISC is symmetry-forbidden, it does not compete with  $S_1 \rightarrow S_0$  IC to recover the ground state. Although ISC between  $S_1$  and  $T_2$  is allowed, our spectroscopic calculations

indicate that the  $T_2$  state of 1-O,O' is  $\sim 1500 \text{ cm}^{-1}$  above  $S_1$ .<sup>10–12</sup> In marked contrast to 1-O,O', the heavy-atom congeners are all emissive, suggesting that their  $T_2$  states lie below  $S_1$ . This supports the idea that fast ISC occurs via a SOC-allowed  $S_1 \rightarrow T_2$  pathway followed by rapid IC to  $T_1$ . The  $S_1 \rightarrow T_2$  ISC is primarily mediated by the degree of Pt d-orbital character admixed into the  $a_2$  HOMO–1,  $b_1$  highest occupied molecular orbital (HOMO), and the  $b_1$  LUMO because  $\langle \text{Pt}(xz) | L_z | \text{Pt}(xy) \rangle \neq 0$ .

For 2-O,S, 3-S,S, 4-S,Se, and 5-Se,Se, three processes are potentially responsible for the observed differences in the  $T_1$  lifetimes (differences in the  $T_1 \rightarrow S_0$  lifetimes may also be influenced by vibrationally-activated nonradiative decay via Pt–E stretching modes; relaxation dominated by these modes should decrease with increased mass of the E donor, but this is not borne out by experiment): (1) the relative  $T_1$ – $S_0$  energy gaps, (2) the orbital allowedness of  $T_1 \rightarrow S_0$  ISC, and (3) ligand SOC. First, although an energy gap law argument might be used to explain why 2-O,S possesses the shortest  $T_1$  lifetime, this does not explain the observed  $T_1$  lifetimes for 3-S,S–5-Se,Se (Table 1) because the  $T_1$  energies ( $E_{\text{PL}}$ ) determined from luminescence measurements are essentially equivalent. Second, ISC between  $T_1$  and  $S_0$  is disfavored in 3-S,S and 5-Se,Se because the group theoretical arguments mentioned above result in  $\langle A_1 | L_z | A_1 \rangle = 0$ . However, bonding calculations indicate a large anisotropic covalency in low-symmetry 2-O,S that manifests itself in a rotation of the Pt( $d\pi$ ) orbital component of the HOMO toward the Pt–S bond, effectively mixing some Pt( $xy$ ) character into the HOMO to give  $\langle \Psi_{T_1} | L_z | \Psi_{S_0} \rangle \propto \langle \text{Pt}(xy) | L_z | \text{Pt}(xz) \rangle \neq 0$  and an enhanced nonradiative ground-state recovery rate. This orbital rotation is markedly less apparent in 4-S,Se because of the similar Pt–S and Pt–Se covalencies (i.e., effective  $C_{2v}$  symmetry), and this is reflected in the similar  $\tau$  values for 3-S,S and 4-S,Se. Finally, ligand SOC represents a potentially important mechanism for modulating  $T_1$  lifetimes across the 3-S,S–5-Se,Se series. Although direct SOC-induced ISC between  $T_1$  and  $S_0$  is orbitally forbidden in  $C_{2v}$  symmetry, mechanisms for overcoming this forbiddenness derive from spin-vibronic and vibronic SOC, which couple states of different spin multiplicity.<sup>13–15</sup> Only out-of-plane  $a_2$  and  $b_1$  vibrations are effective in mixing the out-of-plane  $E_{p_x}$  orbitals with the in-plane  $E_{p_y}$  and  $E_{p_z}$  orbitals. This results in linear combination of atomic orbital (LCAO)-expanded matrix elements of the type  $\langle E_{p_x} | L_z | E'_{p_y} \rangle$  and  $\langle E_{p_x} | L_z | E'_{p_z} \rangle$  being nonzero and a ligand SOC contribution to  $T_1 \rightarrow S_0$  ISC. Because the SOC constant for selenium is  $\sim 4.5$  times that of sulfur ( $\zeta_{\text{Se}} = 1659 \text{ cm}^{-1}$ ;  $\zeta_{\text{S}} = 365 \text{ cm}^{-1}$ ;  $\zeta_{\text{O}} = 154 \text{ cm}^{-1}$ ),<sup>16</sup> one would expect shorter  $T_1$  lifetimes as a function of selenium ligation, and this is observed experimentally across the 3-S,S–5-Se,Se series.

Understanding the electronic structure contributions to  $k_{\text{ISC}}$ ,  $k_r$ , and  $k_{\text{nr}}$  provides keen insight into the continued development of these and related compounds as key components of dyes for solar cells and for photonics. For 1-O,O', the  $T_2$  state lies above

$S_1$  and  $S_1 \rightarrow T_1$  ISC does not compete with rapid IC to recover the ground state. In marked contrast to **1-O,O'**, the close energy match between  $S_1$  and  $T_2$  for the heavy-atom congeners of the series provides a mechanism for enhanced  $S_1 \rightarrow T_2$  ISC. Coupled with energy gap law effects, anisotropic covalency, and ligand SOC, these factors modulate excited-state lifetimes by at least 3 orders of magnitude. Exquisite ligand control of these processes suggests that heteroatom effects can be utilized to further modify excited-state lifetimes in a controlled manner. Efforts along these lines are underway.

## ■ ASSOCIATED CONTENT

### ■ Supporting Information

Synthesis, characterization, and computational details. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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### Notes

The authors declare no competing financial interest.

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